## Bernoulli's Principle

## Introduction

A very important equation in fluid dynamics is the Bernoulli equation. This equation has four variables: velocity $(v)$, elevation $(z)$, pressure $(p)$, and density $(\rho)$. It also has a constant $(g)$, which is the acceleration due to gravity. Here is Bernoulli's equation:
$\frac{v^{2}}{2}+g z+\frac{p}{\rho}=$ constant
To understand and use this equation, we must know about streamlines.

Streamlines are curves that are tangent to the velocity vector of the flow. In other words, they show the direction a fluid element will travel in at any point in time. A streamline is best illustrated by examples:

If you throw a leaf in a stream of water, a streamline is the path of the leaf as it floats downstream. Of course the leaf can take any number of paths depending on where it lands in the stream after it was thrown. Streamlines exist underwater as well. Imagine a submerged particle that is neutrally buoyant (meaning
it neither sinks nor floats). The particle, such as a waterlogged twig, also follows a streamline down the river.

When values for velocity, pressure, etc. are plugged into Bernoulli's equation, the result of the equation will be the same (constant) at every point along the streamline.

Take, for example, a water reservoir at the top of a hill. Imagine a streamline from the reservoir to a residence, where the water is used. At the reservoir, the velocity of the water is very small as it slowly moves toward the spillway. For a water particle near the surface of the reservoir, the pressure is also very small. The elevation above the residence, however, is quite large. Let's plug some of these values into Bernoulli's equation:

$$
\begin{aligned}
& \text { velocity, } y=0.01 \frac{\mathrm{~m}}{\mathrm{~s}} \\
& \text { elevation, } z=1000 \mathrm{~m} \\
& \text { pressure, } p=1 \mathrm{~Pa} \\
& \text { density of water, } \rho=1000 \frac{\mathrm{~kg}}{\mathrm{~m}^{2}} \\
& \text { acceleration due to gravity, } g=9.8 \frac{\mathrm{~m}}{\mathrm{~s}^{2}}
\end{aligned}
$$

## Questions:

1. What is the value of Bernoulli's equation with these values?
2. What are the units of this value?
3. Which of the following has the biggest effect on the value of Bernoulli's equation for this case: velocity, elevation, or pressure?

Now let's follow a particle of water as it leaves the reservoir. As it travels over the spillway and down toward the residence, the elevation decreases, the velocity increases, and if the water travels in an enclosed piping system, the pressure increases. Let's use Bernoulli's equation to find the pressure of the water half-way down toward the residence. We'll assume that we are traveling along the same streamline and that there are no energy losses along the way (more on this later).

## Questions:

4. What is the value of Bernoulli's equation at an elevation of 500 m ?
5. If we measure the speed of the water to be $10 \frac{\mathrm{~m}}{\mathrm{~s}}$ at 500 m of elevation, what is the pressure?

At the end of our journey, the water particle has reached the residence where it is sitting stationary in a faucet, ready to be used by the resident.

## Head Loss

In reality, the actual water pressure at the faucet would be much, much less than 9.8 MPa . The reason for this is due to head loss, which is energy in a moving fluid that is lost due to friction and turbulence in the water as it travels from the reservoir to the residence. Head loss is associated with the length, diameter, and smoothness of the pipe, bends, fittings, and valves.

Turbulence is one fluidic phenomenon that causes head loss. Turbulence is a fluidic region where the particles that make up the fluid are chaotic or random in motion. Turbulence is seen in Figure lin the region behind the wing. In the turbulent zone, the streamlines are continually and quickly changing shape and direction so that they are unrecognizable. Turbulent zones take energy from the fluid and contribute to head loss. Engineers try to minimize turbulence in piping systems in order to reduce the
amount of energy required to move fluid through the piping system. Despite these efforts, all fittings in piping systems such as valves, tees, and unions cause turbulence. Head loss from this turbulence can be estimated for a fitting if the velocity $(v)$ of the water is known. Head loss is also dependent on the type of fitting. The following equation and table can be used to estimate head loss from fittings $\left(h_{f, \text { miner }}\right)$. Note that the acceleration of gravity constant $(g)$ in the following equation.

Head loss, $h_{\text {fminor }}=\frac{F \nu^{2}}{2 g}$
In our previous example, we considered a reservoir at the top of a hill and a piping system that carries the water to residences below. Consider a single leg of the pipeline that carries water from the reservoir to a residence, as illustrated in Figure .


Figure 1. Pipeline from a reservoir to a single residence.

## Questions:

6. What fittings would you choose for this pipeline from the above table to minimize head loss?
7. Assume that the water flowing through the pipe is traveling at $10 \stackrel{m}{s}$. What is the total head loss from the fittings you chose?

Fittings are not the only components in piping systems that cause losses. Friction between the water molecules and the surface of the pipe also contribute to head loss. Factors that influence head loss due to friction are:

- Length of the pipe ( $(t)$
- Effective diameter of the pipe $\left(D_{n}\right)$
- Velocity of the water in the pipe ( $v$ )
- Acceleration of gravity (g)
- Friction from the surface roughness of the pipe ( x )

The head loss due to the pipe is estimated by the following equation:
$h_{f \text { major }}=\lambda \frac{l v^{2}}{2 D_{h g}}$

## Combining Bernoulli's Equation With Head Loss

We can now combine the concepts of Bernoulli's equation and head loss to understand the "Fluidic Energy Equation". We divide each term in Bernoulli's equation by the "specific weight" of the fluid $(\gamma)$. The specific weight of a fluid is simply the density of the fluid multiplied by the acceleration of gravity: $(y=\rho g)$.Thus, the effects of head loss result in:
$\frac{p_{1}}{\gamma}+\frac{v_{1}^{2}}{2 g}+z_{1}=\frac{p_{2}}{\gamma}+\frac{v_{2}^{2}}{2 g}+z_{2}+h_{f, 5 \operatorname{sac} l}$
Recall our discussion of streamlines from the first section. The left side of this equation represents the state of a fluid particle at the start of our streamline, for example sitting in the water reservoir at the top of a hill. The right side of this equation is the state of our fluid particle at the end of our streamline, for example at the water tap of the residence at the bottom of a hill. The energy of the fluid particle in the reservoir is equal to the energy of the fluid particle at the tap plus the head losses. In other words, because of head losses, some of the energy is sapped from the fluid particle during its journey from the reservoir to the tap.

Again, consider our pipeline from the reservoir to the residential tap. Suppose you measure the pressure and velocity at two points in the piping system (along the same streamline) and discover that they are the same for both points. In other words, $v_{1}=v_{2}$ and $p_{1}=p_{2}$.

## DERIVATION OF BERNOULLI'S EQUATION

Bernoulli's equation relates the speed of a fluid at a point, the pressure at that point and the height of that point above a
reference level. It is just the application of work-energy theorem in the case of fluid flow. We shall consider the case of irrotational and steady flow of an incompressible and nonviscous liquid.

Figure below shows such a flow of a liquid in a tube of varying cross-section and varying height. Consider the liquid contained between the cross-sections A and $B$ of the tube. The heights of $A$ and $B$ are h1 and h2 respectively from a reference level.

This liquid advances into the tube and after a time $\Delta t$ is contained between the cross-sections $\mathrm{A}^{\prime}$ and $B^{\prime}$ as shown in

Figure below


Suppose the area of cross-section at $A=A 1$
the area of cross-section at $B=A 2$
the speed of the liquid at $A=v 1$
the speed of the liquid at $B=v 2$
the pressure at $A=$ the pressure at $B=P 2$
and the density of the liquid $=\dot{\rho}$.
The distance
$A A^{\prime}=v l \Delta t$
and the distance
$B B^{\prime}=v 2 \Delta t$.
The volume between $A$ and $A^{\prime}$ is
A 1 vldt and
the volume between B and $B^{\prime}$ is
A2v2 $\Delta t$.
By the equation of continuity,
$A 1 v 1 \Delta t=A 2 v 2 \Delta t$
The mass of this volume of liquid is
$\Delta m=p A l v 1 \Delta t=p A 2 v 2 \Delta t$.
Let us calculate the total work done on the part of the liquid just considered. The forces acting on this part of the liquid are
(a) P 1A1, by the liquid on the left
(b) P2 A2, by the liquid on the right
(c) $(\Delta m) g$, the weight of the liquid considered and
(d) cell, contact forces by the walls of the tube.

In time $\Delta t$, the point of application of P 1 Al is displaced by $A A^{\prime}=v l \Delta t$.

Thus, the work done by P1A1 in time $\Delta t$ is
$\mathrm{W} 1=(\mathrm{P} 1 \mathrm{~A} 1)(v 1 \Delta t)=\mathrm{P} 1[\Delta \mathrm{~m} / \rho]$

Similarly, the work done by P2 A2 in time $\Delta \mathrm{t}$ is
$\mathrm{W} 2=-(\mathrm{P} 2 \mathrm{~A} 2)(\mathrm{v} 2 \mathrm{At})=-\mathrm{P} 2[\mathrm{~m} / \mathrm{\rho}]$

The work done by the weight is equal to the negative of the change in gravitational potential energy.

The change in potential energy (P.E.) in time $\Delta t$ is
P. E. of $A^{\prime} B B^{\prime}-P$. E. of A A'B
$=$ P. E. of $A^{\prime} B+$ P. E. of BB' - P. E. of AA' - P. E. of A'B
$=P$. E. of BB' P. E. of $\mathrm{AA}^{\prime}$
$=(\Delta \mathrm{m}) \mathrm{gh} 2-(\Delta \mathrm{m})$ gh1 .
Thus, the work done by the weight in time $\Delta t$ is
$\mathrm{W} 3=(\Delta \mathrm{m}) \mathrm{gh} 1-(\Delta \mathrm{m}) \mathrm{gh} 2$.

The contact force $\dot{\mathrm{N}}$ ( does no work on the liquid because it is perpendicular to the velocity.

The total work done on the liquid considered, in the time interval At, is

$$
\begin{align*}
\mathrm{W}=\mathrm{W} 1+\mathrm{W} 2 & +\mathrm{W} 3 \ldots \ldots \ldots \ldots \ldots(1)  \tag{1}\\
\mathrm{W} 1 & =(\mathrm{P} 1 \mathrm{~A} 1)(\mathrm{v} 1 \Delta \mathrm{t})=\mathrm{P} 1[\Delta \mathrm{~m} / \dot{\rho}] \\
\mathrm{W} 2 & =-(\mathrm{P} 2 \mathrm{~A} 2)(\mathrm{v} 2 \mathrm{At})=-\mathrm{P} 2[\Delta \mathrm{~m} / \dot{\rho}] \\
\mathrm{W} 3 & =(\Delta \mathrm{m}) \mathrm{gh} 1-(\Delta \mathrm{m}) \mathrm{gh} 2 .
\end{align*}
$$

Therefore,

$$
\begin{aligned}
\mathrm{W} & =\mathrm{W} 1+\mathrm{W} 2+\mathrm{W} 3 \\
& =\mathrm{P} 1[\Delta \mathrm{~m} / \dot{\rho}]-\mathrm{P} 2[\Delta \mathrm{~m} / \dot{\rho}]+(\Delta \mathrm{m}) \mathrm{gh} 1-(\Delta \mathrm{m}) \mathrm{gh} 2 .
\end{aligned}
$$

Workdone on the system is equal to the change in its kinetic energy. Thus,
$\mathrm{P} 1[\Delta \mathrm{~m} / \mathrm{\rho}]-\mathrm{P} 2[\Delta \mathrm{~m} / \dot{\rho}]+(\Delta \mathrm{m}) \mathrm{gh} 1-(\Delta \mathrm{m}) \mathrm{gh} 2$
$=1 / 2[\Delta \mathrm{~m}] \mathrm{v} 2^{2}-1 / 2[\Delta \mathrm{~m}] \mathrm{v} 1^{2}$
$\mathrm{P} 1 / \dot{\rho}+\mathrm{gh} 1+(1 / 2) \mathrm{v} 1^{2}=\mathrm{P} 2 / \dot{\rho}+\mathrm{gh} 2+(1 / 2) \mathrm{v} 2^{2}$
Or
$\mathrm{P} 1+\dot{\rho} \mathrm{gh} 1+(1 / 2) \dot{\rho} \mathrm{v} 1^{2}=\mathrm{P} 2+\dot{\rho} \mathrm{gh} 2+(1 / 2) \dot{\rho} \mathrm{v} 2^{2}$

## Or

$\mathrm{P}+\dot{\rho} \mathrm{gh}+(1 / 2) \dot{\rho} \mathrm{v}^{2}=$ constant
This is known as Bernoulli's equation.

